

[0011] Magnetic field region **18** is a band of asymmetric conductance region **12** subjected to a localized magnetic field **B** oriented normal to the plane of thermoelectric structure **10** and out of the page, as depicted in FIG. 1. This magnetic field may be provided by depositing and polarizing magnetic material **28**, a thin layer of magnetic material located atop and/or beneath the plane of thermoelectric structure **10**, adjacent magnetic field region **18**. Alternatively, magnetic field **B** may be an external field from, e.g., a magnetized coil. Magnetic field **B** is herein assumed for simplicity to be substantially uniform within magnetic field region **18**, although all implementations of magnetic field **B** will of course vary somewhat over magnetic field region **18**. Collimating regions **20** and **22** are bands of asymmetric conductance region **12** situated between magnetic field region **18** and shorting bars **14** and **16**, respectively. Collimating regions **20** and **22** are regions with negligible magnetic field that serve to collimate ballistic charge flow between shorting bars **14** and **16** (in either direction).

[0012] Collimating regions **20** and **22** feature collimating guides **24**, while magnetic field region **18** features curved guides **26**. Collimating guides **24** and curved guides **26** are physical discontinuities along lines in thermoelectric structure **10** that act as scattering barriers to form channels in collimating regions **20** and **22** and magnetic field region **18**, respectively, by ballistically scattering incident charge carriers. Adjacent collimating guides **24** may, for instance, be separated by a distance of approximately 1 nm—approximately 1 μm along an axis parallel to shorting bars **14** and **16**, depending on the material and operating temperature of thermoelectric structure **10**. Adjacent curved guides are separated by a similar distance. Collimating guides **24** and curved guides **26** may be created in a variety of ways, including by laser or mechanical scribing, surface level doping, field doping, or lithographic patterning. Collimating guides **24** are straight, parallel lines that act to focus charge carrier trajectories in collimating regions **20** and **22** along transport direction **T** or the opposite direction, $-T$. Curved guides **26** focus charge carriers moving in transit direction **T** from shorting bar **14** to shorting bar **16**, but act to continually frustrate charge transport in opposite direction $-T$, as described in further detail below. Collimating guides **24** and curved guides **26** extend throughout the entire thickness of thermoelectric structure **10**.

[0013] It is well known from elementary physics that a charge carrier of charge q , when travelling with vector velocity v through a magnetic field characterized by vector B , will experience a Lorentz force:

$$F = qv \times B. \quad [\text{Equation 4}]$$

[0014] A charge travelling in a plane through a magnetic field normal to that plane thus experiences a Lorentz force qvB in the plane and at right angles with v . The direction of curvature of a charge trajectory due to Lorentz force is opposite for conductors travelling with velocities v and $-v$, and of opposite signs q and $-q$. As depicted in FIG. 1, an electron travelling in transport direction **T** will deflect to the left under magnetic field **B**, while an electron travelling in the opposite direction $-T$ will deflect to the right under magnetic field **B**. Curved guides **26** take advantage of this broken symmetry by allowing substantially unobstructed electron flow in transit direction **T** while frustrating electron flow in the opposite direction $-T$. Curved guides **26** form parallel curved channels in magnetic field region **18** that coincide with the arcs of

curvature of forward conduction (i.e. in transit direction **T**), and thus more closely match the natural deflection trajectories of negative charge carriers moving in transit direction **T** than in the opposite direction $-T$. Thus, electrons travelling in transport direction **T** scatter on curved guides **26** substantially less and at wider angles than electrons travelling in the opposite direction $-T$. This asymmetry results in longer ballistic trajectories in the $-T$ direction than in transmit direction **T**, with corresponding forward electrical conductivity $\sigma_{\text{forward}} > \text{reverse electrical conductivity } \sigma_{\text{reverse}}$. This behavior is illustrated and described in further detail with respect to FIGS. 2A and 2B.

[0015] FIGS. 2A and 2B depict ballistic trajectories of negative charge carriers such as electrons through magnetic field region **18**. FIG. 2A shows the trajectory of a charge carrier moving in transport direction **T**, while FIG. 2B shows the trajectory of a charge carrier moving in opposite direction $-T$. In both cases the Lorentz force causes the charge carrier to deflect in a counter-clockwise direction, according to the right-hand rule. In FIG. 2A, the charge carrier is deflected substantially to the right along a path defined by curved guides **26**, and scatters at large angles with respect to curved guide **26**. This scattering adds relatively little to the total path length of the charge carrier trajectory in FIG. 2A, corresponding to a high value of forward electrical conductivity σ_{forward} . In FIG. 2B, by contrast, the charge carrier is deflected substantially to the right, and scatters several times at progressively smaller angles with respect to curved guides **26**. This scattering dramatically lengthens the total path length of the charge carrier trajectory in FIG. 2B, corresponding to a low value of reverse electrical conductivity $\sigma_{\text{reverse}} < \sigma_{\text{forward}}$.

[0016] Thermoelectric structure **10** enables high values of ZT . Magnetic field regions **18** with perpendicularly applied magnetic fields **B** and curved guides **26** coinciding with arcs of curvature of charge carriers traveling in transport direction **T** can yield ratios of forward to reverse conductance $\sigma_{\text{forward}} / \sigma_{\text{reverse}} \sim 10$, potentially enabling the creation of $ZT > 5$ thermoelectric materials which would fundamentally change the coefficient of performance of solid state materials, potentially opening up their use for all solid state commercial refrigeration systems.

[0017] While the invention has been described with reference to an exemplary embodiment(s), it will be understood by those skilled in the art that various changes may be made and equivalents may be substituted for elements thereof without departing from the scope of the invention. In addition, many modifications may be made to adapt a particular situation or material to the teachings of the invention without departing from the essential scope thereof. Therefore, it is intended that the invention not be limited to the particular embodiment(s) disclosed, but that the invention will include all embodiments falling within the scope of the appended claims.

The status of the claims is as follows:

1. A thermoelectric structure comprising:

- a thin thermoelectric film extending in a plane between parallel first and second shorting bars; and
- a plurality of curved ballistic scattering guides formed in a magnetic field region of the thin thermoelectric film subjected to a local, substantially uniform, nonzero magnetic field normal to the plane of the thin thermoelectric film.

2. The thermoelectric structure of claim 1, wherein the shape of the curved ballistic scattering guides substantially matches an arc of curvature of a charge carrier travelling in a